

1     **On the Benefits of a High-Resolution Analysis for Convective**  
2     **Data Assimilation of Radar Observations using a Local Ensemble**  
3                     **Kalman Filter**

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## ABSTRACT

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6 The Deutscher Wetterdienst (DWD) is developing an implementation of the Local Ensemble  
7 Transform Kalman Filter (LETKF) for the cloud resolving COSMO model. This study  
8 shows in an idealized convective testbed that the LETKF is able to perform storm-scale  
9 Data Assimilation of simulated Doppler radar observations.

10 Perfect model experiments are used to investigate how the limited predictability of con-  
11 vective storms affects precipitation forecasts by comparing a fine scheme with low analysis  
12 error to a coarse scheme that allows variance regarding position, shape and occurrence of  
13 storms in the ensemble. To get there, the coarse scheme uses averaged superobservations  
14 and a coarser evaluation of the analysis weights, a larger localization radius and a weaker  
15 Gaussian constraint on the analysis solution. Performing 3-hour forecasts of convective sys-  
16 tems with typical lifetimes exceeding 6 hours, forecasts from the detailed analyses of the  
17 fine scheme are found to be advantageous to those of the coarse scheme during the first 1-2  
18 hours, regarding the predicted storm positions. After 3 hours in the convective regime used  
19 here, the forecast quality of the different schemes appears indiscernible, judging by RMSE  
20 and verification methods for rain-fields and objects.

21 It is concluded that, for operational assimilation systems, the analysis might not neces-  
22 sarily need to be detailed on the grid scale of the model. Depending on the forecast lead  
23 time, and on the presence of orographic or synoptic forcings that enhance the predictability  
24 of storm occurrences, analyses from a coarser scheme might suffice. As a positive side-effect,  
25 the computational cost of the Kalman Filter solution can be reduced strongly.

# 1. Introduction

In the last decade, Data Assimilation (DA) of Doppler radar observations using an Ensemble Kalman Filter (EnKF) (Evensen 1994; Houtekamer and Mitchell 1998) has been explored to be a feasible method to obtain suitable initial states of convective storms for very short-range ensemble forecasts. As Stensrud et al. (2009) note, a goal of convective DA is to be competitive to nowcast-warning systems that have generally been superior to model forecasts without data assimilation for at least the first three hours of leadtime (Kober et al. 2012).

Previous studies that used radar data in an EnKF focused on storm-analyses that apply relatively small scales of  $\sim 10$  km for the covariance localization and converge the analysis ensemble closely towards the observations, pursuing the goal of obtaining initial states with low errors as the basis of their ensemble forecasts. This study tries to assess the performance of that converging approach, as it is costly and may be of limited value due to the short predictability of convection in the chaotic atmospheric system where small errors grow rapidly (Lorenz 1969; Lilly 1990). The framework of the Local Ensemble Transform Kalman Filter (LETKF) (Hunt et al. 2007) is used here in an idealized setup of the nonhydrostatic COSMO model (Baldauf et al. 2011) with severe and long lived convection.

In contrast to the converging approach, the alternative of a coarse assimilation scheme is proposed which is cheaper to compute and focuses less on the small-scale convergence of the analysis ensemble; variance and errors of the analysis should be allowed that are larger in scale *and* amplitude. For lead times of a few hours, after the strongly nonlinear growth of small errors has saturated (Zhang et al. 2003), this coarse scheme may not be subject to much penalty in forecast accuracy in comparison to a fine and converged scheme. Secondly, the coarse scheme with its larger variance may avoid dynamical imbalances caused by analysis increments than can arise when the member states of the localized analyses are rigorously converged (Greybush et al. 2011).

The sections of the introduction below are devoted to the problems of converging analyses

53 and limited predictability and how they could be adressed using assimilation schemes of  
54 different precision. Finally, the setup for the numerical experiments is laid out.

55 *a. Converged storm analyses in previous studies*

56 The main observation types in convective EnKF are reflectivity and radial Doppler  
57 wind. All studies on this topic mentioned here used the EnSRF-algorithm (Houtekamer  
58 and Mitchell 2001) or a similar algorithm that localizes the background error covariance  
59 matrix  $\mathbf{P}^b$ , provided by a background ensemble of 50-100 members, to increase the effec-  
60 tive sample size. In their first OSSE study on convective EnKF, Snyder and Zhang (2003)  
61 identified spurious storms in the members to be detrimental to the analysis solution. Tong  
62 and Xue (2005) were able to suppress spurious storms by assimilating observations of no-  
63 reflectivity outside the observed storms; their positive results were confirmed in the real data  
64 cases of Aksoy et al. (2009, 2010). As the ensemble mean is the solution that minimizes the  
65 cost function of the Kalman Filter, most studies focused on the improvement of this mean  
66 solution. The non-Gaussian and multimodal distribution of the observed dynamical systems  
67 (Dance 2004) forces the probability distribution of the resulting analyses, constrained to be  
68 Gaussian by the filter, to have a deficient amount of variance. A common solution for this  
69 is the inflation of the background error covariance, preferably adaptive to the respective  
70 situation (Anderson 2008). For real observations, Dowell and Wicker (2009) and Dawson  
71 et al. (2012) added noise during the assimilation inside the observed storms and Aksoy et al.  
72 (2009) perturbed the vertical wind profile. Both methods resulted in better analyses with  
73 a more consistent variance. Wang et al. (2012) applied the “Running in Place”-method of  
74 Kalnay and Yang (2010) on convective EnKF experiments, resulting in a better spin-up of  
75 the storms and a generally lower error level.

76 Using a Gaussian-like correlation function for the localization (Gaspari and Cohn 1999),  
77 all studies agreed on the usage of a small localization cutoff radius between 4 and 12 km,  
78 depending on the spatial stage of the storm development (Sobash and Stensrud 2013) and

79 assumed an observation error of approximately 5 dBZ for (no-)reflectivity and  $1 \text{ m s}^{-1}$  for  
80 wind observations in the observation error covariance matrix  $\mathbf{R}$ . The resulting precise storm-  
81 analyses without spurious convection are hereby referred to as “converged” or “fine” analysis  
82 ensembles because all members essentially agree on position and shape of the observed  
83 storms.

84 *b. Limited predictability of convective storms*

85 Lilly (1990) and Skamarock (2004) estimated the predictability of mesoscale convective  
86 systems to be in the range of tens of minutes to 1 hour before the upscale error growth (Lorenz  
87 1969) taints the forecast completely. Zhang et al. (2003), Hohenegger and Schär (2007)  
88 and Done et al. (2012) compared randomly perturbed forecasts of organized convection to  
89 unperturbed reference runs. They found small-scale perturbations to grow very quickly  
90 and nonlinearly, reaching a growth-saturation at 3-6 hours. The specific predictability limit  
91 in these studies depended on the presence of a large scale forcing that predetermined the  
92 occurrence of convection spatially and modally. Craig et al. (2012) used perturbations from  
93 a Latent Heat Nudging assimilation scheme for convective storm forecasts and concluded  
94 a lower predictability for convection in regimes without synoptic forcing where the storms  
95 positions are mostly random.

96 *c. Fine and coarse assimilation schemes*

97 Taking the limited predictability of convection into account, and the fact that forecast  
98 products of local models are usually disseminated with lead times  $> 3 \text{ h}$ , the benefits of the  
99 fine and converging approach of the previous studies on convective EnKF may not justify  
100 the cost: *A coarse and less converged analysis ensemble could provide equally good forecasts*  
101 because small perturbations grow quicker than large-scale disturbances (Lorenz 1969). As  
102 in Kuhl et al. (2007), the deviations of the analysis members from the reference run are

103 regarded here as the perturbations which grow during the forecast.

104 In this paper, a fine and a coarse assimilation scheme are described. The aim of the fine  
105 scheme is to reproduce the converging ensemble behavior of the previous studies mentioned  
106 in Section 1a. The fine analysis localizes the covariances on a scale of a few kilometers and  
107 uses high resolution observations. It will therefore contain observed details on the smallest  
108 resolved scales of the model grid – which is smaller than the “effective resolution” of the  
109 model (Skamarock 2004) where the numerical solution of the equations can be considered  
110 accurate. The LETKF-algorithm used here localizes in the observation space, so it will be of  
111 interest if the results are comparable to those studies that localize the background covariance  
112 in model space. In contrast, the coarse scheme is configured to have perturbations that are  
113 larger with respect to scale, amplitude and horizontal positions of the updraft cores and are  
114 thus subject to a slower error growth (Lorenz 1969). It is based on the fine scheme and uses a  
115 scaling factor of 4 for most of the assimilation parameters: (i) a larger horizontal localization  
116 radius is applied, (ii) coarsened superobservations are used that provide information only  
117 on this larger scale and (iii) the local analysis weights of the LETKF are computed on  
118 a coarse grid and then spatially interpolated. The latter method deprives the analysis of  
119 some precision but allows the linear combination of the analysis members to vary more  
120 smoothly (Yang et al. 2009). The experiments in both the fine and the coarse scheme will  
121 show that it is also necessary to (iv) adjust  $\mathbf{R}$  by inflating it in order to attenuate the  
122 rigorous convergence towards the observations, as the application of an untreated  $\mathbf{R}$  results  
123 in noisy and imbalanced analysis member states that exhibit poor physical resemblance to  
124 the reference run.

125 Applying (i)-(iv) should allow states in the coarse analysis that are more consistent  
126 with the model physics because larger horizontal portions of background member storms  
127 are contained in the analysis member storms. As a technical benefit, the lower number of  
128 observations and the lower analysis resolution result in a significant computational accelera-  
129 tion of the analysis. Attributing the superobservations with less weight in the analysis using

130 method (iv) will push the members less rigorously towards the observations by analysis in-  
131 crements (cf. Fig 1). This aims for a multimodal posterior distribution containing variance  
132 not only storm-internally but also with respect to the occurrence of storms themselves, with  
133 the eventuality of spurious convection. These properties define the “non-converged”, coarse  
134 assimilation scheme.

135 In addition to the convergence properties of the two experiments, other issues that are  
136 caused by the localized analysis of non-observed variables, such as a cold bias of the tem-  
137 perature variable, are explored.

#### 138 *d. Numerical experiments*

139 The experiments are performed in a perfect model environment to focus solely on the  
140 predictability of the dynamics, excluding effects of model error. Simulated radar observations  
141 are drawn from a nature run that contains organized convective systems with lifetimes over  
142 six hours. A cycled assimilation covers a time-span of three hours, followed by three hours of  
143 ensemble forecast. For all experiments, the EnKF will be challenged by a convective ensemble  
144 that is spun up from random initial noise and therefore lacks any prior knowledge about  
145 the position of the observed storms in the nature run. This “bad background” makes the  
146 convergence of the ensemble challenging for the filter, but it is free from possibly beneficial  
147 influence e.g. of convective triggers that are introduced outside the EnKF algorithm at  
148 predetermined locations (Aksoy et al. 2009, 2010).

149 The 3-hour forecasts from both schemes are evaluated using the RMSE of the states  
150 together with object- and field-based forecast skill scores to evaluate how quickly the advan-  
151 tage of a fine analysis state is lost in the forecast with respect to the coarse analysis, due to  
152 the limited predictability.

## 2. Fine and coarse assimilation scheme

This section first describes the data assimilation setup, consisting of nature run, synthetic observations, convective ensemble and LETKF-algorithm, followed by the implementation of the fine and the coarse scheme. The Kilometer-scale ENsemble Data Assimilation (KENDA) system (Reich et al. 2011) is being developed at the Deutscher Wetterdienst (DWD). It couples an LETKF-implementation with an ensemble of the COSMO-DE model. The latter is used for operational forecasting over Germany (Baldauf et al. 2011).

COSMO solves the full non-hydrostatic and compressible Navier-Stokes equations using a time-splitting Runge Kutta approach for fast and slow tendencies in the prognostic wind variables  $U$ ,  $V$  and  $W$  and the deviations of temperature  $T'$  and pressure  $p'$  from a stationary hydrostatic base state. The moist physics use six state variables of water vapor, cloud water, cloud ice, rain, snow and graupel with a first order bulk microphysics scheme. A radiation scheme for long- and shortwave radiation is applied. Surface fluxes of latent and sensible heat are parametrized and constrained by constant surface temperature and surface specific humidity throughout the simulation.

### *a. Nature run*

This study uses the testbed setup of COSMO with idealized initial state, periodic boundary conditions (BCs) and a homogeneous flat landscape as the lower boundary. A convection-permitting horizontal resolution of 2 km and 50 vertical levels are used in a domain of  $396 \times 396 \times 20$  km extent. The initial profile of all model runs is horizontally homogeneous and based on the sounding of Payerne (CH, Radiosonde 06610) at 12 UTC on Juli 30th 2007, a day with strong convective storms and mesoscale convective systems (MCS), favored by a high CAPE value of  $2200 \text{ J kg}^{-1}$  together with a vertical wind shear that allows organized convection with heavy precipitation and propagating gust fronts (Bischof 2011).

Instead of initializing convection with predefined warm bubbles (Aksoy et al. 2009) or



178 targeted noise (Dowell et al. 2004; Tong and Xue 2005) with amplitudes that directly trigger  
179 thermals, uncorrelated grid point noise is added at the initial time  $t_0$  to the temperature field  
180  $T$  and the vertical wind speed  $W$  in the boundary layer (BL) with amplitudes of 0.02 K and  
181  $0.02 \text{ m s}^{-1}$ . The model runs start at 06 UTC and rapidly develops a convective boundary  
182 layer. Triggering instabilities are generated through the effect of the radiation scheme on  
183 the tropospheric temperature.

184 First preliminary showers evolve in random locations at 08 UTC which grow until 10  
185 UTC and mostly die off by 12 UTC (Fig. 2). The surviving systems grow into intense  
186 storms and MCS by 14 UTC and propagate with the mean wind in a  $45^\circ$ -direction through  
187 the domain with lifetimes  $\geq 6 \text{ h}$  (Fig. 3). The horizontally contiguous rain areas extend  
188 between 30 and 200 km with a reflectivity larger than 30 dBZ. Surface fluxes of sensible and  
189 latent heat allow the cold pools to relax in the wake of the storms. The periodic BC allow  
190 the storms to grow in a way that is “natural” for the model physics in the given sounding.  
191 The time-window between 14 and 20 UTC is chosen for 3 hours of cycled assimilation until  
192 17 UTC, followed by ensemble forecasts with 3 hours leadtime until 20 UTC.

193 Five random realizations of the initial noise in the BL are used as nature runs to cover  
194 the possible storm characteristics and positions.

### 195 *b. Synthetic observations*

196 Synthetic radar observations of reflectivity and radial windspeed are generated for each  
197 realization. The observations are computed on a regular grid with a horizontal spacing  
198 of 2 km and a vertical spacing of 1 km, covering the heights between 500 m and 12500  
199 m, in order to mimic a region of good radar coverage. Reflectivity  $Z$  is computed from  
200 mixing ratios of graupel ( $QG$ ), rain ( $QR$ ) and snow ( $QS$ ) using the simple implementation  
201 of Done et al. (2004). As in other OSSE studies (Tong and Xue 2005), Gaussian noise  
202 with a standard deviation of  $\sigma_{refl} = 5 \text{ dBZ}$  is added to simulate errors of measurement and  
203 representativity. The reflectivity observations  $Z$  are masked to regions where  $Z > 5 \text{ dBZ}$ .

204 Below this threshold, they are regarded as observations of no-reflectivity (Tong and Xue  
 205 2005; Aksoy et al. 2009) with a nominal value of 0 dBZ and  $\sigma_{no-refl} = 2.5$  dBZ, assuming  
 206 a good post-processing that can identify rain-free areas. The horizontal wind component  $U$ ,  
 207 with an added error of  $\sigma_U = 1$  m s<sup>-1</sup>, is used as an analogue for radial wind observations.  
 208 This would be an accurate approximation if the storms are far away from the radar site. The  
 209 regular observation geometry and the usage of  $U$  ensures that the data coverage is uniform  
 210 and all storms are equally well observed.

211 *c. Initial ensemble*

212 The synthetic radar observations are assimilated using an ensemble of  $k = 50$  members  
 213 which differ from the nature run and among themselves only in the random distribution of  
 214 the initial noise. The spin-up time between 06 and 14 UTC enables the members to contain  
 215 storms with similar characteristics but completely uncorrelated horizontal positions. This  
 216 approach was chosen to deprive the ensemble of any prior knowledge about the reference  
 217 when the assimilation starts, as it could have been provided by a “manual” positioning of  
 218 warm bubbles in the members or a confinement of the initial noise to regions of observed  
 219 reflectivity (Tong and Xue 2005; Dowell and Wicker 2009).

220 *d. Implementation of the LETKF*

221 To produce an analysis ensemble, the LETKF-algorithm (for a full description see the  
 222 original paper) determines the  $\mathbf{w} = \bar{\mathbf{w}}^a$  that minimizes the cost function

$$J^*(\mathbf{w}) = (k - 1)\mathbf{w}^T\mathbf{w} + [\mathbf{y}^o - \bar{\mathbf{y}}^b - \mathbf{Y}^b\mathbf{w}]^T\mathbf{R}^{-1}[\mathbf{y}^o - \bar{\mathbf{y}}^b - \mathbf{Y}^b\mathbf{w}] \quad (1)$$

223 where the  $k$ -dimensional vector  $\mathbf{w}$  defines the optimal linear combination of ensemble  
 224 member states that minimizes  $J^*$ .  $\mathbf{y}^o$  contains the observations and  $\mathbf{R}$  is the observation  
 225 error covariance matrix.  $\bar{\mathbf{y}}^b$  and  $\mathbf{Y}^b\mathbf{w}$  are given by approximating the observation operator  
 226  $H$  to be linear about the  $m$ -dimensional background ensemble mean state  $\bar{\mathbf{x}}^b$

$$H(\bar{\mathbf{x}}^b + \mathbf{X}^b \mathbf{w}) \approx \bar{\mathbf{y}}^b + \mathbf{Y}^b \mathbf{w} \quad (2)$$

227 where  $\mathbf{X}^b$  is a  $m \times k$  matrix whose columns are given by the deviations of the single  
 228 forecast members from their mean  $\mathbf{x}^{b(i)} - \bar{\mathbf{x}}^b$  and

$$\mathbf{Y}^b = H(\mathbf{x}^{b(i)}) - k^{-1} \sum_{i=1}^k H(\mathbf{x}^{b(i)}) \quad (3)$$

229 The analysis of (1) are computed locally for every analysis grid point. However, they do  
 230 not necessarily need to be taken out on the full model resolution: the spatial field of the local  
 231  $\mathbf{w}^{a(i)}$  is usually quite smooth, so a coarser “analysis grid” can be chosen horizontally and  
 232 vertically on which the local analysis weights  $\mathbf{w}^{a(i)}$  are first determined and then interpolated  
 233 onto the model grid (Yang et al. 2009).

234 For the local analysis, only nearby observations are taken into account by localizing the  
 235 observation error covariance matrix  $\mathbf{R}$  with a Gaussian-like correlation function (Gaspari  
 236 and Cohn 1999) that is zero where the distance  $r$  of the single observations is larger than  
 237 the “cutoff-length”  $r_{Loc}$  of the localization radius. The dimension of the local  $\mathbf{R}$  corresponds  
 238 then to the number of local observations.

239 Since periodic boundary conditions have not been implemented in the LETKF, observa-  
 240 tions closer to the border of the domain than  $r_{Loc}$  are discarded. The untreated  $\mathbf{R}$  contains  
 241 the  $\sigma^2$ -variances mentioned in Section 2b and is diagonal in the present case because the  
 242 added errors are uncorrelated.

243  $\mathbf{R}$  here is the only manipulable factor in (1) that determines how closely the analysis  
 244 members are drawn towards the observations by the filter. As the localized assimilation of  
 245 observations produced by nonlinear operators can lead to imbalances (Greybush et al. 2011),  
 246 it appears sensible to inflate  $\mathbf{R}$  to lessen the impact of the observations in (1), as will be  
 247 shown later.

249 The fine analysis scheme R4 (Table 1) is based on previous convective-scale data assimi-  
 250 lation studies (Section 1a). It is devised as a control run to reproduce previous results and  
 251 to serve as a benchmark for the performance of the coarse experiment later.

252 A horizontal localization length scale  $l_{Loc,h} = 4$  km is chosen that corresponds to a Gaus-  
 253 sian correlation function with  $G(4 \text{ km}) = \exp(-1/2) \approx 0.61$ . A Gaussian-like correlation  
 254 function  $C$  is used that goes to zero within the finite distance of the *localization radius*  
 255  $r_{Loc,h} = 2l_{Loc,h}\sqrt{10/3} \approx 14.6$  km (Gaspari and Cohn 1999; Hamill et al. 2001). This  $l_{Loc,h}$  is  
 256 chosen for all observation types, as previous studies found it to be an effective compromise  
 257 (Sobash and Stensrud 2013). The resolution of the synthetic observations is  $\Delta x_{obs} = 2$  km,  
 258 same as the model resolution. The vertical localization length scale,  $l_{Loc,v}$ , ranges from 3  
 259 km near the surface to 5 km at the model top. An assimilation interval  $\Delta t_{ass}$  of 5 minutes  
 260 represents the typical availability of volume observations from a scanning Doppler radar (Lu  
 261 and Xu 2009). The analysis grid here has the full horizontal resolution of  $\Delta x_{ana} = 2$  km.  
 262 The vertical resolution, provided by the 20 levels  $n_{lev}^{ana}$  of the analysis grid, is lower than the  
 263 model resolution of 50 levels but should suffice, given the strong overlap due to the vertical  
 264  $l_{Loc,v}$ . The weights  $\mathbf{w}^{a(i)}$  are multiplied by a constant covariance inflation factor of  $\rho = 1.05$   
 265 in order to enhance the span of the analysis state space.

266 In this R4-setup, the positions of the analysis storms should coincide with the observed  
 267 storms. Within the storm-cores, the analysis states are expected to be detailed with a low  
 268 error and small variance, while spurious convection is suppressed outside of them by assim-  
 269 lating volume-observations of no-reflectivity. These are the requirements on a converged and  
 270 detailed analysis ensemble formulated in the introduction.

271 As it turned out in preliminary experiments, providing the filter with the untreated  $\mathbf{R}$   
 272 can result in imbalanced increments. These taint the solution by introducing gravity wave  
 273 noise that arises from dynamically inconsistent model states when the members are drawn  
 274 too closely to the observations. It was therefore chosen to inflate the entries of  $\mathbf{R}$  from

275  $(\sigma_U = 1 \text{ m/s})^2$  to  $(\sigma_U = 5 \text{ m/s})^2$ , from  $(\sigma_{refl} = 5 \text{ dBZ})^2$  to  $(\sigma_{refl} = 20 \text{ dBZ})^2$  and from  
 276  $(\sigma_{no-refl} = 2.5 \text{ dBZ})^2$  to  $(\sigma_{no-refl} = 20 \text{ dBZ})^2$  (cf. Table 1). Lessening the influence of the  
 277 observations in (1) by this method resulted in more consistent and stable analyses as with  
 278 the untreated **R**.

279 The deteriorating effect appeared to be most critical with the reflectivity observations  
 280 that are computed by a very nonlinear operator whose increments in model space are  
 281 nonetheless approximated to be linear by (2). To test this hypothesis, R4\_forced was devised  
 282 as a sensitivity experiment wherein **R** contains the reflectivity error values with which the  
 283 synthetic observations were originally perturbed:  $(\sigma_{refl} = 5 \text{ dBZ})^2$  and  $(\sigma_{no-refl} = 2.5 \text{ dBZ})^2$   
 284 (cf. Table 1).

#### 285 *f. Coarse scheme R16*

286 In the coarse scheme R16, more horizontal variance should be allowed, so a larger  $l_{Loc,h} =$   
 287 16 km is chosen and the observational resolution  $\Delta x_{obs}$  is decreased from 2 km to 8 km. At  
 288 the same time, the resolution of the analysis grid  $\Delta x_{ana}$  is decreased from 2 km to 8 km,  
 289 so there are  $4^2$  times fewer local analysis vectors  $\mathbf{w}^{a(i)}$  to be computed by the local cost  
 290 functions (1), reducing the computational effort and smoothing the analysis field further.  
 291 Consequently, for every local analysis the same number of observations is used in R16 as in  
 292 – but now on a coarser scale.

293 The observations are coarse-grained by horizontally averaging the values and positions  
 294 of  $4 \times 4$  blocks of the original observations into one central superobservation (SO); this is  
 295 preferable to data-thinning and reduces the information to the desired coarse scale (Alpert  
 296 and Kumar 2007; Salonen et al. 2009; Seko et al. 2004).

297 The analysis increments should therefore consist of larger parts of storms that are more  
 298 dynamically consistent internally than in R4, as the analysis weights of the different members  
 299 will vary less between adjacent model grid points and the influence radius of the observations  
 300 is larger. In one single analysis, a whole storm-core with up- and downdraft could be drawn

301 from the background ensemble in R16, whereas in R4 the updraft-region of an analysis storm  
 302 might originate from a different background member than the downdraft region.

303 The fact that the coarse and large-scale scheme R16 has the same number of observa-  
 304 tions per local analysis as the fine scheme R4 suggests that the absolute influence of the  
 305 observations on the local solution of (1) should also be the same. As this is regulated by the  
 306 magnitude of the local  $\mathbf{R}^{-1}$ -entries, R16 uses an inflated  $\mathbf{R}$ -matrix that contains the error  
 307 same values as in R4 (cf. Table 1), but now for the SOs.

308 This reasoning is tested by the sensitivity-experiment R16\_forced which uses the error  
 309 values for the SOs which have been effectively lowered by the averaging: The original  
 310 observations of R4 are a normally distributed random variable (cf. Section 2b), so the  
 311 block-averaging in R16 should reduce the error standard deviation of the SO by a factor of  
 312  $\sqrt{1/4^2} = 1/4$  to multiply  $\mathbf{R}$  with. If this is done, the observational part of the local cost  
 313 function (1) becomes larger, because the local analysis has the same number of observations  
 314 but now with a larger  $\mathbf{R}^{-1}$ . This rigorous formulation is tested in R16\_forced whose  $\mathbf{R}$ -  
 315 entries are derived from the control experiment R4. The entries of  $\mathbf{R}_{R16\_forced}$  with SOs are  
 316  $(\sigma_{U,SO} = 1.25\text{m/s})^2$ ,  $(\sigma_{refl,SO} = 5 \text{ dBZ})^2$  and  $(\sigma_{no,refl,SO} = 5 \text{ dBZ})^2$  (cf. Table 1).

317 The the inflated  $\mathbf{R}$ -matrix of R16 is intended to allow a not-converged ensemble wherein  
 318 even the occurrence of an observed storm, given by observations of  $U$  and reflectivity, should  
 319 have a variance; also, spurious convection should be permitted by lessening the suppressive  
 320 influence of no-reflectivity observations. These are the requirements on a non-converged and  
 321 coarse analysis ensemble formulated in Section 1.

322 The observations in R16 are assimilated using a longer cycling interval  $\Delta t_{ass}$  of 20 minutes  
 323 between subsequent analyses. In sensitivity experiments for R16 with  $\Delta t_{ass} = 5$  minutes,  
 324 it was found that the inevitable noise of the increments was introduced too frequently and  
 325 deteriorated the analysis. Also, significant large-scale differences in the members storm  
 326 structures could not spin up during  $\Delta t_{ass} = 5$ , rendering the analyses worse than with  
 327  $\Delta t_{ass} = 20$  min. The resulting noisy structures were similar to the disturbances of R16\_forced

328 that are shown later, albeit with lower amplitude.

329 Taking advantage of the reduced memory requirements of the coarser horizontal grid,  
 330  $n_{lev}^{ana} = 25$  had been chosen for R16, giving a modest improvement over  $n_{lev}^{ana} = 20$ . The  
 331 slight advantage over R4 with respect to the additional vertical levels is assumed to be offset  
 332 by the larger number of discarded observations in the wider border regions of  $l_{Loc,h} = 16$  km,  
 333 so R16 and R4 should generally be comparable.

334 *g. RMSE, spread and consistency ratio*

335 The accuracy of analysis and forecast states  $\mathbf{x}$  is measured by the Root Mean Square  
 336 Error (RMSE), computed in model space for the different model variables:

$$RMSE(\mathbf{x}) = \sqrt{m^{-1} \sum_{l=1}^m (x_l^{nature} - x_l)^2} \quad (4)$$

337 where  $m$  is the number of grid points. The corresponding variance is the spread  $\overline{spr}$  of the  
 338 ensemble  $\mathbf{x}^i$  around its mean  $\bar{\mathbf{x}} = k^{-1} \sum_{i=1}^k \mathbf{x}^i$ , given by

$$\overline{spr} = m^{-1} \sum_{l=1}^m \sqrt{(k-1)^{-1} \sum_{i=1}^k (x_l^i - \bar{x}_l)^2} \quad (5)$$

339 To fulfill the Gaussian assumption of the filter, the ensemble spread should represent the  
 340 actual error of the analysis, so the consistency ratio  $CR$

$$CR = \frac{\overline{spr}}{RMSE} \quad (6)$$

341 should be  $CR = 1$ .

342 Unlike studies that only regarded convectively active regions above a certain reflectivity  
 343 threshold, the present study evaluates RMSE and spread for all grid points of the domain.  
 344 This results a) in a generally low error level because many small error values of clear air  
 345 regions contribute to the mean error and b) in analysis errors which should be regarded  
 346 relatively to the free error level of an ensemble running in parallel without assimilation, as  
 347 the overall variance of the convective situation may change during the diurnal cycle.

### 3. Assimilation results at fine and coarse resolution

First the fine experiment R4 is evaluated during the assimilation window 14-17 UTC in Section 3a, then the results of R16 are presented for comparison in Section 3b. An example of the results for one realization of the nature run is given in Fig. 4, which compares snapshots of nature run 01 to the analysis means of R4 and R16 at 17 UTC. The maximum column reflectivity is the vertical 2D-projection of the 3D-reflectivity field.  $T$  at the height of  $z = 150$  m shows the cold pools,  $W$  at  $z = 3500$  m shows regions of up- and downdrafts. Fig. 5 shows the RMSE and spread of the ensemble means during the assimilation window and during the forecast window from 17 to 20 UTC for the model variables  $U$ ,  $W$ ,  $T$  and  $QR$ . Analogously, Fig. 6 shows the mean RMSE of the single members. The dash-dotted lines in the RMSE-Figures provide an error level by showing the error of a free-running ensemble that has not undergone any assimilation.

#### a. Performance of R4

The fine scheme R4 is able to converge the ensemble mean onto the observed storms for the directly observed precipitation variables contained in the reflectivity observations, so the reflectivity field of the R4-mean closely resembles the nature run by 17 UTC (Fig. 4). The RMSE of the rain mixing ratio  $QR$  is strongly reduced in the analysis mean of R4 (Fig. 5) and exhibits a very low spread: the graupel mixing ratio ( $QG$ ) behaves analogously (not shown). This points at a strong convergence of the ensemble members onto the observed clouds. Fig. 7 compares the reflectivity of some representative analysis-members of R4 to the nature run. All member-storms look similar to the nature run and spurious convection is largely suppressed by observations of no-reflectivity. In Fig. 9, the analysis ensemble distributions are evaluated at selected points of updrafts where the nature run shows a vertical velocity of  $+10 \text{ m s}^{-1}$  (a) or a reflectivity of 40 dBZ (b) at 17 UTC. These are locations where observations are favorably present: the reflectivity is observed directly, the vertical velocity



373 of the analysis is provided through background covariances where reflectivity observations  
374 and horizontal convergence observations ( $U$ ) indicate an updraft. The analysis member  
375 states of R4 have converged closely around these values with a distribution that exhibits a  
376 Gaussian shape.

377 Judging by Figs. 4, 7 and 9, the cycled analyses produce up- and downdrafts – although  
378 only provided by background error covariances – in the right positions, shapes and amplitudes  
379 (Fig. 4); the RMSE of  $W$  is clearly reduced for the R4-mean relative to the free-running  
380 ensemble, and a good consistency ratio is reached by maintaining ensemble spread (Fig. 5).

381 The RMSE of the horizontally smooth  $T$ -field is not reduced as much relative to the free  
382 error level as for the horizontally intermittent variables  $W$  and  $QR$ . The analysis increments  
383 even appear to be largely detrimental in the case of R4 (showing as a reversed sawtooth  
384 pattern), especially when the single members are considered (Fig. 6). Fig. 4 also shows that  
385 the shapes and positions of the analysis cold pools are in accordance with the nature run  
386 but exhibit a cold bias. This bias is assumed to be caused by the covariance localization:  
387 the mean temperature is not conserved when increments add local linear combinations of  
388 members, (Janjić et al. 2012, 2013) possibly resulting in convective systems that are generally  
389 too wet and therefore develop too strong cold pools due to evaporative cooling. The deficient  
390  $T$ -analysis resulting from EnKF-assimilation of radar data has been noticed in previous  
391 studies (Zhang et al. 2004; Dong et al. 2011). The previous studies found that temperature  
392 observations from a dense surface mesonet can reduce the problem of too strong cold pools,  
393 but this method is not considered here.

394 Observations of the horizontally smooth  $U$ -field are available where reflectivity is present.  
395 The filter acts beneficially on  $U$  for R4 and reduces the RMSE relative to the free error level  
396 and also absolutely. As with  $T$ , the consistency ratio of  $U$  exceeds 1 by 17 UTC, indicating  
397 that the chosen covariance inflation factor  $\rho = 1.05$  might be too large for these variables.  
398 The  $V$ -field is closely correlated to  $U$  by the dynamics of the storms and the  $45^\circ$  background  
399 wind field and behaves analogously to  $U$  in terms of errors and spread.

400 Fig. 10 compares R4 to the sensitivity experiment R4\_forced which has an uninflated  $\mathbf{R}$ -  
401 matrix for (no-)reflectivity observations. For all variables, the errors are larger in R4\_forced,  
402 and the detrimental effect on  $T$  shows even more clearly. In (2) it was assumed that a  
403 deviation from the mean in the observation space is directly proportional to the attributed  
404 deviation in the model space. The inflated  $\mathbf{R}$  in R4 allows a greater weight for the background  
405 ensemble in (1) and thus does not follow the assumption of (2) as rigorously which then acts  
406 beneficially on the quality of the analysis.

407 Summarizing these results, the R4-scheme appears to produce analyses that are com-  
408 parable to previous convective EnKF studies with radar data. The covariance localization  
409 in observation space of the LETKF-algorithm appears to work well, compared to the lo-  
410 calization in model space in the EnSRF-filters of other studies. The storms in the analysis  
411 ensemble resemble the storms of the nature run closely in terms of position, amplitude and  
412 dynamics. *The mean of the strongly converged ensemble is representative for the best so-*  
413 *lution, the single analysis members show a strong agreement with little variance inside the*  
414 *observed storms.* These were the requirements for a fine and converged scheme formulated  
415 in the introduction. R4 is regarded to fulfill them.

#### 416 *b. Performance of R16*

417 Now the assimilation results of the coarse scheme R16 are evaluated. With the horizon-  
418 tally smooth variable  $U$ , Fig. 5 exhibits a similar error for the R16-mean as for the R4-mean;  
419 the slight advantage of R4 for  $U$  appears even less significant if the mean of the member-  
420 RMSEs in Fig. 6 is considered where no difference between the two schemes is apparent.  
421 The analysis quality of  $T$  seems to be better for R16 – the analysis increments are always  
422 beneficial here (Fig. 6). R16 assembles the analysis storms from horizontally larger portions  
423 of member-storms, apparently reducing the dynamical inconsistencies in smooth variables  $U$   
424 and  $T$  that are introduced by the localization. Nevertheless, also R16 exhibits a cold bias in  
425 the ensemble mean, especially due to additional cold pools of spurious storms.

426 For the horizontally intermittent rain-variables and thus for the reflectivity, the analysis  
 427 quality of R16 is inferior to R4. In Fig. 4, the counterparts of the observed storms in the  
 428 R16-mean have roughly the right position and shape but are blurred, caused by variability  
 429 of the updraft core positions within the ensemble (Fig. 8). This was expected due to the  
 430 combination of the SOs that provide information only on a coarser scale, a larger  $l_{Loc,h}$  and  
 431 coarser analysis grid that facilitate horizontally smoother increments, and the inflated  $\mathbf{R}$   
 432 that allows a higher level of variance. The resulting R16-ensemble is not as converged as the  
 433 R4-ensemble: its distribution is broader, less Gaussian and even multimodal in locations of  
 434 updrafts and strong reflectivity in the nature run (Fig. 9). Spurious storms exist largely and  
 435 show up in the reflectivity fields of the ensemble members (Fig. 8) and the mean (Fig. 4).  
 436 The RMSE-values of the variables related to vertical motion ( $W$  and  $QR$ ) are larger for R16,  
 437 lying just slightly below the free error level (Fig. 6), but now possess significant spread in  
 438  $QR$  (Fig. 5).

439 The introduction of the coarse scheme required a larger variance for the storm positions  
 440 than in the fine scheme, together with larger analysis error and more spread, allowing also  
 441 spurious convection. These requirements are regarded to be fulfilled by the coarse scheme  
 442 R16.

443 Fig. 11 compares a typical storm system of one nature run to the corresponding analysis  
 444 means of R16 and the sensitivity experiment R16\_forced, the latter with an uninflated  $\mathbf{R}$ -  
 445 matrix for (no-)reflectivity observations. In contrast to R16, the mean  $W$ -field of R16\_forced  
 446 reveals single updraft cores with large amplitudes where all members agree on their positions,  
 447 recognizable also in the larger amplitude of the mean reflectivity field: *R16\_forced corresponds*  
 448 *to a stringently converged version of R16, but unlike to R4, the converged mean does not*  
 449 *resemble the physical properties of the nature run.* The strong and imbalanced increments of  
 450 R16\_forced cause gravity waves that show in the  $W$ -field where noise is present in the analysis  
 451 mean. Strong updrafts lie close to downdrafts in an alternating pattern which appears to be  
 452 dynamically inconsistent with the nature run.

453 The effective sample size in the R16-experiments is lower than in the R4-experiments due  
454 to the larger localization radius (Hunt et al. 2007). Producing a converged ensemble with  
455 horizontally broad increments as in R16\_forced would possibly need a larger ensemble size –  
456 with the given 50 members, the non-converged ensemble of R16 turned out to work better.

## 457 **4. Ensemble forecasts from fine and coarse resolution** 458 **analyses**

459 From the last analysis ensembles of the experiments at 17 UTC, ensemble forecasts are  
460 performed until 20 UTC.

### 461 *a. RMSE forecast error*

462 For  $W$  and  $QR$ , the RMSE in R4 converges to the error level of R16 within 1-2 hours,  
463 both for the ensemble mean (Fig. 5) and the members (Fig. 6). Within this period, R4  
464 quickly regains spread for  $QR$ , caused by the precipitating cores of the member storms that  
465 are moving apart from the well-determined position of the converged analysis. At 20 UTC,  
466 the degree of horizontal blurring in R4 equals that of R16 (Fig. 12).

467 For  $U$  and  $T$ , the RMSE of the forecasts grows more slowly with no significant difference  
468 between R4 and R16. Qualitative differences of the analysis perturbations in R4 and R16  
469 with respect to the reference run are thus contained mostly in the horizontally intermittent  
470 variables such as  $W$  and  $QR$ .

### 471 *b. Field-based evaluation by the DAS-score*

472 Anticipating that much of the error in the R4 and R16 ensemble mean forecast fields  
473 (Fig. 12), two spatial verification measures are chosen to compare the forecast fields to the  
474 nature run. Fields of maximum column reflectivity are chosen for this evaluation because

475 high reflectivity is usually accompanied with strong precipitation and winds – these two are  
476 the essential threats to be predicted by a convective storm forecast. The observation field  
477 of the nature run here is thresholded above 10 dBZ to separate the storms and eventually  
478 overlapping anvils.

479 The Displacement and Amplitude Score DAS (Keil and Craig 2009) uses a pyramidal  
480 matching algorithm to compare two fields by an optical flow technique. A vector field is  
481 computed that morphs the forecast onto the reference field and vice versa, using a maxi-  
482 mum search radius of 45 km here. The average magnitude of the displacement vector field,  
483 normalized by maximum search radius, defines the DIS-component of the DAS-score, and is  
484 displayed in Fig. 13. During the assimilation window, R4 is clearly superior to R16 due to  
485 the converged ensemble. This advantage is lost during the forecast window within 1 hour –  
486 after that, the forecast quality of storm positions is indistinguishable between the fine and  
487 coarse schemes.

488 *c. Object-based evaluation by the SAL-score*

489 The SAL-score (Wernli et al. 2008) compares statistical properties of thresholded rain-  
490 objects. The structure or  $S$ -component indicates if the forecast objects are smoother and  
491 broader than the observations ( $S > 0$ ) or spikier ( $S < 0$ ), with  $S = 0$  indicating the  
492 correct structure. Only average object properties are compared, without matching, so  $S$  is  
493 independent of location errors and biases. In Fig. 14a, R4 and R16 appear qualitatively  
494 equal with  $S \approx 0$  at 17 UTC, meaning that the right type of objects are present in the  
495 members after the assimilation cycling; the observed convective modes are reproduced, as  
496 far as the projection of the column is concerned. This property holds during the forecast for  
497 the members (dashed lines), but for the ensemble mean (solid lines), the blurring of forecast  
498 features is apparent with an increasing  $S > 0$ .

499 The amplitude or  $A$ -component (Fig. 14b) compares the domain-wide total reflectivity  
500 between forecast and observation and thus displays the overall bias. During the assimilation

501 window, R4 and R16 behave differently. In the first hour of R4, most spurious storms are  
502 subsequently suppressed while the inserted storms have not yet fully developed, resulting  
503 in a negative bias with  $A < 0$ , that then overcompensates to a positive bias with  $A > 0$   
504 by 17 UTC. In R16, the analysis increments always increase the bias by inserting observed  
505 storm-increments without complete suppression of spurious storms, counteracting the relax-  
506 ing tendency of the dynamics during the forecast intervals. During the first 90 minutes of  
507 the forecast window, the  $A$ -bias of R16 decreases, while small and incompletely suppressed  
508 spurious storms in R4 grow quickly and enhance the R4-bias, but eventually converge with  
509 R16. After 2 hours of forecast, R4 and R16 are indistinguishable.

510 The  $L$ -component (Fig. 14c) determines the location error by measuring the horizontal  
511 deviation of the centroid of the forecast reflectivity field from the centroid of the observations.  
512 The results here are similar to those of the DIS-score: R4 has an advantage over R16 by the  
513 end of the analysis period, but this is lost within 1-2 hours of forecast time.

514 The curves of the location error  $L$  are less continuous than the DIS-score and sometimes  
515 even decrease; this is a consequence of the small sample five random realizations that are  
516 evaluated here. Comparing the  $L$ -error between the experiments and the free error level here  
517 is more meaningful than the absolute  $L$ -values.

518 The results of DAS-DIS and the SAL score were verified by evaluating the Brier Score  
519 (Wilks 2006) of the forecast probabilities of R4 and R16 with respect to convective events  
520 exceeding a threshold of 10 dBZ (Fig. 15). A similar convergence in skill can be seen, and  
521 at 20 UTC there is no significant difference of skill between R4 and R16 in the sample of  
522 the five repetitions of the experiments.

## 523 **5. Summary and Discussion**

524 The aim of this study on convective scale data assimilation was to assess the benefits of  
525 high-resolution analysis schemes in the provision of initial conditions for 3-hour precipitation

526 forecasts, taking into account the limited predictability of convective systems due to scale-  
527 dependent error growth.

528 Data Assimilation of long-lived and organized convective systems was performed under  
529 the assumption of a perfect model, using a LETKF in an idealized testbed with simulated  
530 Doppler radar observations of reflectivity and radial wind. The storms were triggered ran-  
531 domly in the convection-permitting COSMO-model using radiative forcing and initial small-  
532 amplitude random noise in a horizontally homogeneous environment with periodic boundary  
533 conditions.

534 Updraft positions, storm intensities and cold pool structures were determined well by the  
535 cycled analyses of the high-resolution assimilation scheme R4; this appeared as a successful  
536 reproduction of previous studies on this topic. The usability of the LETKF-algorithm,  
537 which localizes in observation space, for detailed storm-analyses with radar data is therefore  
538 demonstrated. Reducing the influence of observations on the analysis increments by an  
539 inflated error covariance matrix  $\mathbf{R}$  proved useful for obtaining dynamically consistent model  
540 states.

541 The coarse scheme R16 was devised to produce less precise analyses than R4 by increas-  
542 ing the horizontal localization length scale from 4 to 16 km, by reducing the resolution of  
543 observations and of the analysis grid from 2 to 8 km and by inflating  $\mathbf{R}$  further; by these  
544 means, variance was introduced into the ensemble with respect to the position and occurrence  
545 of updrafts. Spurious convection was permitted in multimodal posterior distributions.

546 Both schemes showed a cold bias due to the introduction of localized analysis increments  
547 that were too wet and caused excessive cooling in low levels.

548 The 3-hour period of cycled assimilation of the two schemes was followed by ensemble  
549 forecasts with a lead time of 3 hours. During the first hour, precipitation forecasts from  
550 analyses of the fine scheme were clearly advantageous to analyses of the coarse scheme in  
551 terms of the storm positions and internal structure, measured by RMSE and the object based  
552 scores DAS and SAL. This advantage of the fine analyses was lost within the first 1-2 hours

553 of forecast by the rapid error growth of the small perturbations.

554 *a. Conclusions*

555 These results suggest that, due to the limited predictability, convective forecasts with  
556 typical leadtimes of 3 hours might not benefit from storm analyses that are detailed on the  
557 scale of the model grid; a coarser representation of storm positions and occurrence as in the  
558 coarse scheme might be sufficient. The computational acceleration hereby is notable.

559 Nonetheless, a fine analysis with low errors is superior for the first hour and could be  
560 useful for forecasts of phenomena smaller than the storm system such as tornados and gust  
561 fronts.

562 *b. Discussion*

563 Looking back at the methods, the coarse scheme R16 is based on the fine scheme R4  
564 which reproduced previous results from other studies. The means by which the coarsen-  
565 ing of the analysis was achieved in R16, namely the larger localization together with a less  
566 detailed analysis computation and coarsened observations, were only useful in combination  
567 with an inflated  $\mathbf{R}$ -matrix to relax the Gaussian constraint of the cost function (shown by  
568 R16\_forced compared to R16). The  $\mathbf{R}$ -inflation was chosen empirically to let the system  
569 produce reasonable analyses. More sophisticated methods exist for the estimation of appro-  
570 priate observation errors for  $\mathbf{R}$  (Desroziers et al. 2005) which are also implemented in the  
571 KENDA-system but were not applied here. Evaluating the effects of the coarsening methods  
572 of R16 separately could give much insight but would be beyond the scope of this paper.

573 The cold bias of the analyses hints at the substantial problem that mass and energy is not  
574 conserved in analyses produced by a localized EnKF (Janjić et al. 2012, 2013). Covariance  
575 localization in observation space as done here shows no substantial differences to localization  
576 in model space.



577 The forecast results displayed the limited predictability of the dynamics in large con-  
578 vective systems. In an operational model, the predictability will be enhanced by effects of  
579 synoptic and orographic forcing (Hohenegger and Schär 2007) that might overlay the efforts  
580 of a fine analysis; in that case, the larger analysis errors of a coarse scheme might be even less  
581 disadvantageous than in the idealized situation presented in this study. Effects of model  
582 error could further diminish the benefits of the fine scheme, especially if it tries to converge  
583 the model states into convective modes that are not supported by the model physics and  
584 the predicted sounding (Stensrud and Gao 2010) – here the gentler approach of the coarse  
585 scheme could be profitable for dynamic consistence.

586 The spatial and temporal quantities of dynamics and predictability that are discussed  
587 here depend on the model resolution, the model physics and the type of long-lived storm  
588 that is simulated. A resolution of 2 km is not sufficient to simulate storm-internal variance  
589 on the scale of single plumes which come addressable with resolutions of 250 m and less  
590 (Bryan and Morisson 2012). Using such a model that is able to resolve three-dimensional  
591 turbulence, an even finer assimilation scheme could be applied to further investigate the  
592 limits of predictability.

593 In the case of a different sounding and the resulting different convective modes such  
594 as mesocyclones, multicell storms and linear squall lines, the predictability limit will be  
595 influenced by the degree of storm-internal organization (Aksoy et al. 2010). The quantitative  
596 time limit found in this study can thus not be generalized for all kinds of convection.

597 The lagged detection of storm cells by radar observations remains a problem. Further  
598 studies could investigate if an ensemble that contains spurious storms might be advantageous  
599 when an assimilation system has to catch storms that develop in locations where clear air  
600 was observed and assimilated before: it may be harmful to the dynamics to select single  
601 precipitating forecast members that have to form large increments for the analysis mean  
602 (Lien et al. 2013).

603 Assimilating reflectivity observations is still difficult due to their non-Gaussian distribu-

604 tions, making the suppression of spurious convection insufficient because (a) the background  
605 distribution of convection is typically intermittent and therefore non-Gaussian and (b) there  
606 is no “negative rain” to assimilate in unaltered observations (Craig and Würsch 2012). A  
607 Gaussian transformation of precipitation observations could help here to fulfill the Gaussian  
608 assumptions of the EnKF (Lien et al. 2013).

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TABLE 1. Experiments as described in Section 2. R4\_f and R16\_f are R4\_forced and R16\_forced.  $l_{Loc,h}$  is the horizontal localization lengthscale,  $r_{Loc,h}$  is the cutoff length of the horizontal covariance localization function,  $\Delta x_{obs/ana}$  is the horizontal resolution of observations and of the analysis grid,  $n_{lev}^{ana}$  is the number of vertical levels of the analysis grid,  $\Delta t_{ass}$  is time-interval between the two consequent analyses,  $\sigma_U$ ,  $\sigma_{refl}$  and  $\sigma_{no-refl}$  are the standard deviations of the observations (for R16 and R16\_f: SOs) that are contained in the **R** for the respective experiment. Some **R**-entries are inflated, see text.

	<b>R4</b>	<b>R16</b>	<b>R4_f</b>	<b>R16_f</b>	
$l_{Loc,h}$	4	16	4	16	km
$r_{Loc,h}$	14.6	58	14.6	58	km
$\Delta x_{obs/ana}$	2	8	2	8	km
$n_{lev}^{ana}$	20	25	20	25	
$\Delta t_{ass}$	5	20	5	20	min
$\sigma_U$	5	5	5	1.25	m s <sup>-1</sup>
$\sigma_{refl}$	20	20	5	5	dBZ
$\sigma_{no-refl}$	20	20	2.5	5	dBZ

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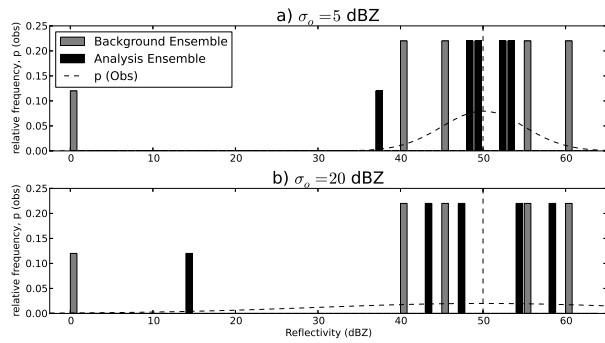


FIG. 1. Relative frequencies of predicted reflectivities for a relatively good background ensemble (gray) and two analysis-ensembles a) and b) (black), computed by the LETKF-algorithm. The observed value is 50 dBZ. In a) a small observation error of  $\sigma = 5$  dBZ is reported to the filter via  $\mathbf{R}$ , in b) an inflated value of  $\sigma = 20$  dBZ. The dashed curve is the assumed normal distribution around the observation.

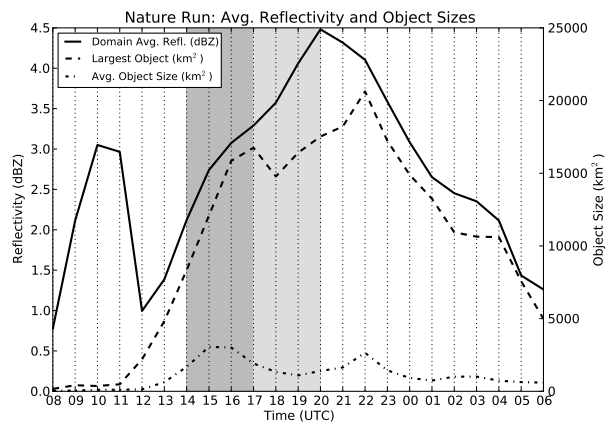


FIG. 2. Time-series of the nature run (averaged over five realizations): Domain-average of the maximum column reflectivity in dBZ (solid), together with average horizontal (dash-dotted) and maximum (dashed) size of the rain-objects, thresholded to  $> 5$  dBZ. Assimilation window is between 14 and 17 UTC (shaded gray), forecast window is between 17 and 20 UTC (shaded pink). Peaks of the object-sizes can be due to merger of anvils of separate convective systems.

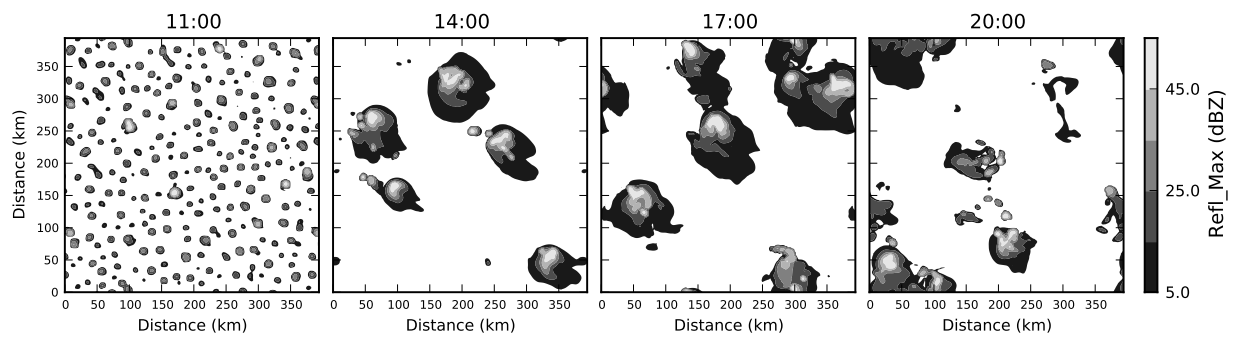


FIG. 3. Time series of maximum column reflectivity of nature run 01. Note that the decline in average reflectivity from 17 to 20 UTC is not representative for the five repetitions (cf. Fig. 2)

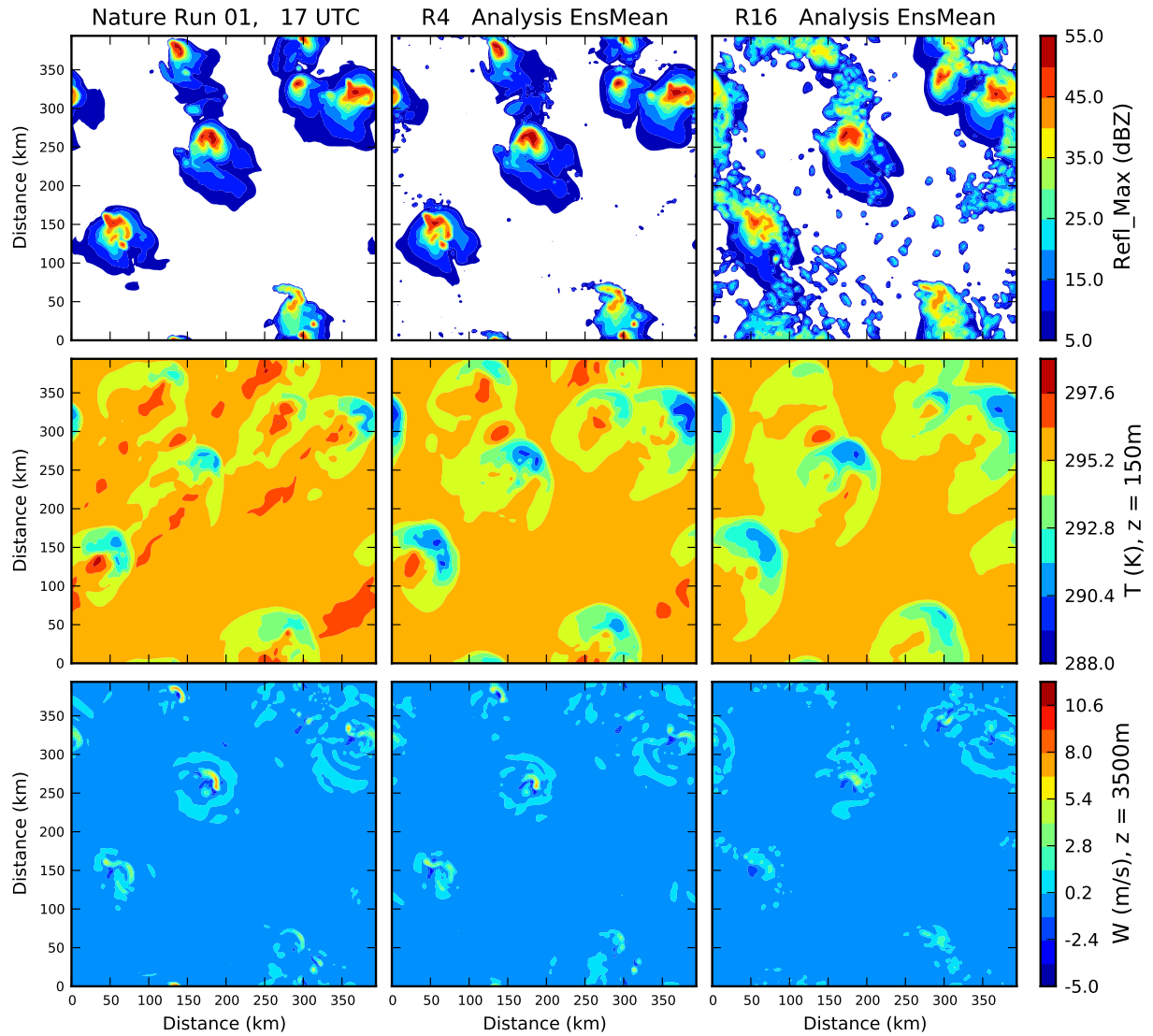


FIG. 4. Nature Run and Analysis Means of R4, and R16 (both of repetition R01) at 17 UTC. Top rows: Maximum reflectivity. Middle rows: Temperature  $T$  at  $z = 150$  m. Bottom rows: Vertical velocity  $W$  at  $z = 3500$  m.



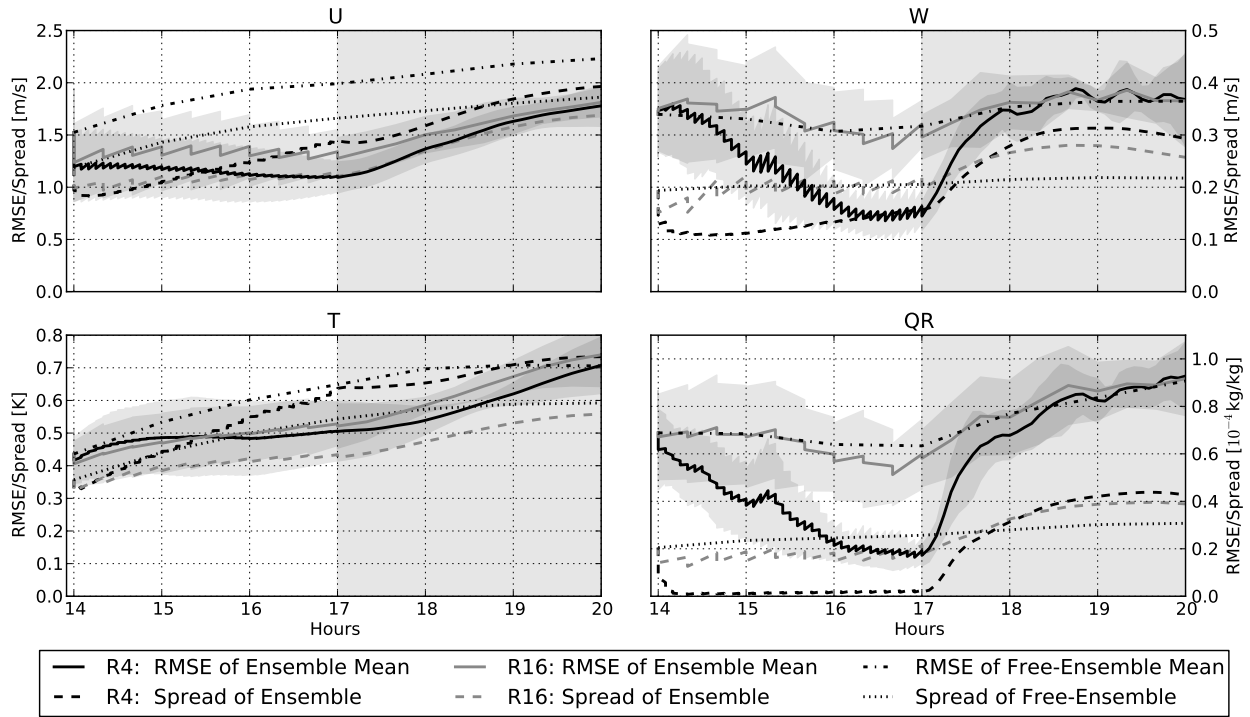


FIG. 5. RMSE and spread of the ensemble means of R4 (black) and R16 (gray) during assimilation (14-17 UTC, forecast ensemble mean and analysis ensemble mean as one consecutive line) and ensemble forecast (17-20 UTC, only forecast ensemble mean). The gray lines show the RMSE and spread of the mean of the free ensemble without any assimilation. All gridpoints are evaluated. The error values are averaged over five realizations of the experiment; an envelope of one empirical standard deviation is depicted by the shaded areas.

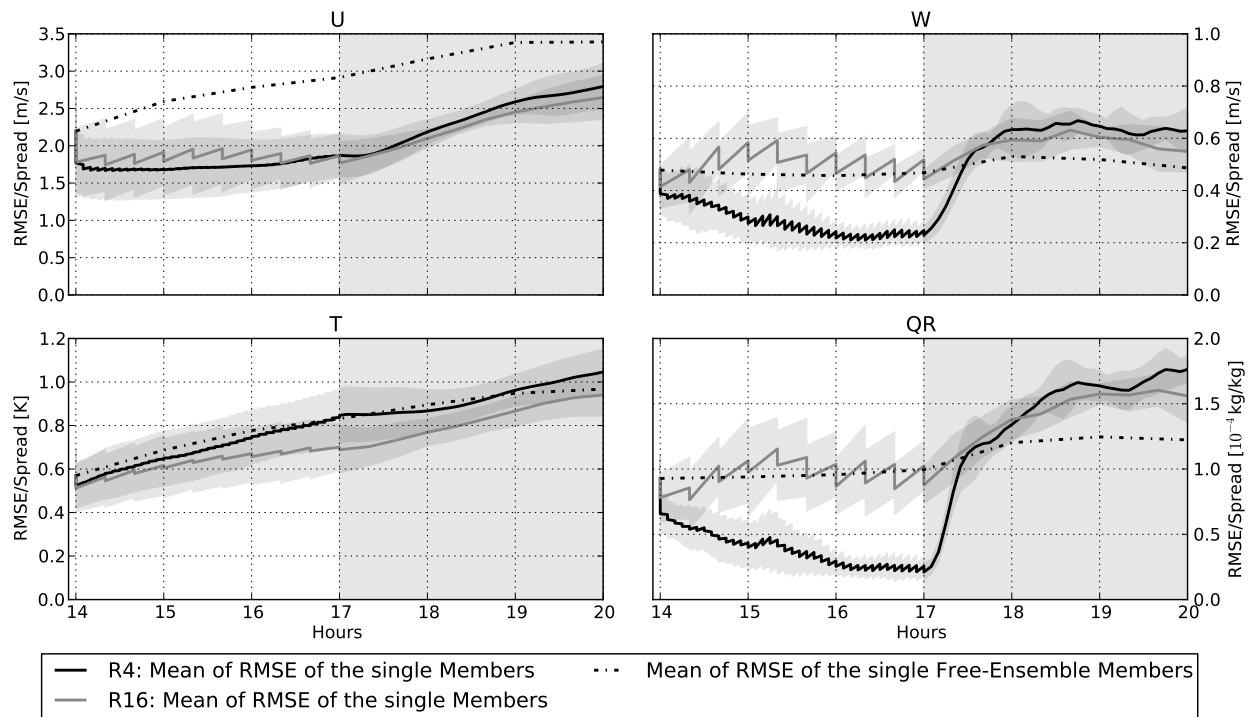


FIG. 6. Like Fig. 5, but for the mean of the RMSE of the single ensemble members of R4 (black) and R16 (gray) instead of the RMSE of the ensemble-mean states.

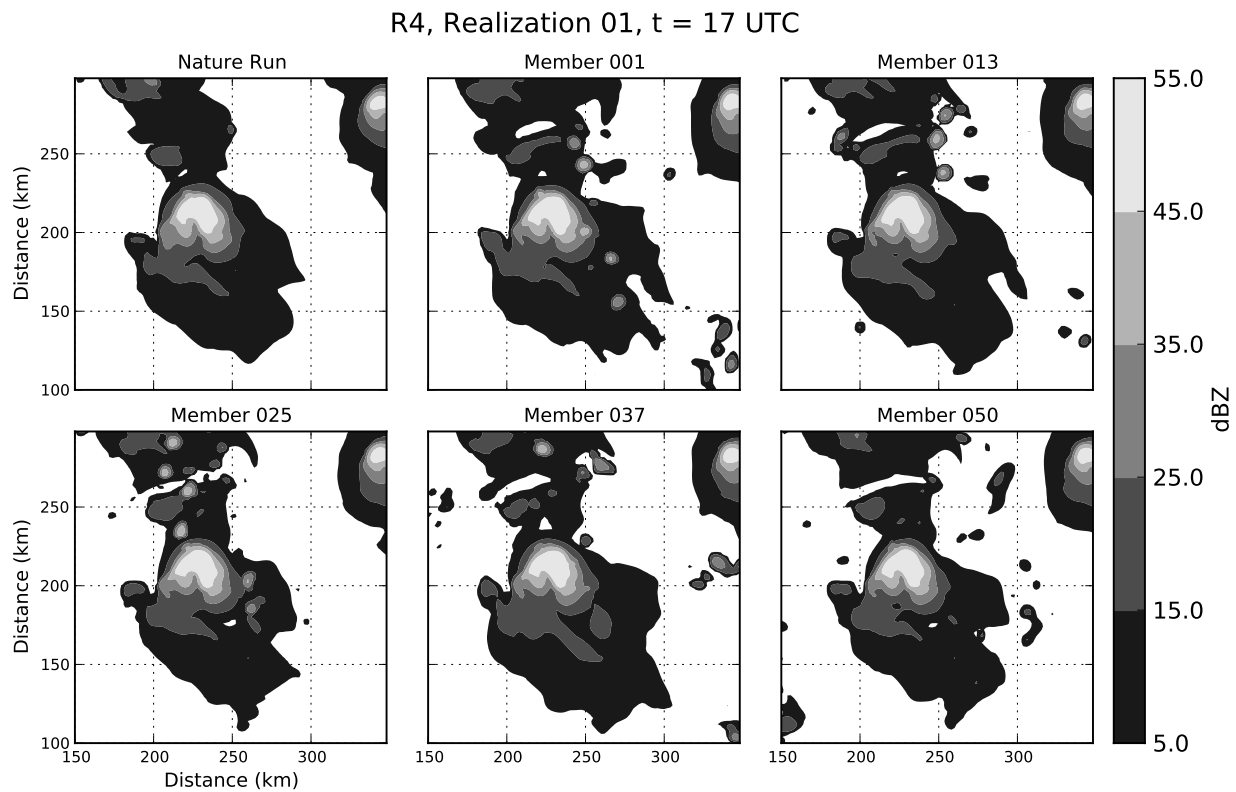


FIG. 7. Maximum reflectivity ( $REFL\_MAX$ ) of nature run and analysis ensemble members 1,13,25,37,50 of R4 (Realization R01) at the last assimilation time 17:00.

R16, Realization 01, t = 17 UTC

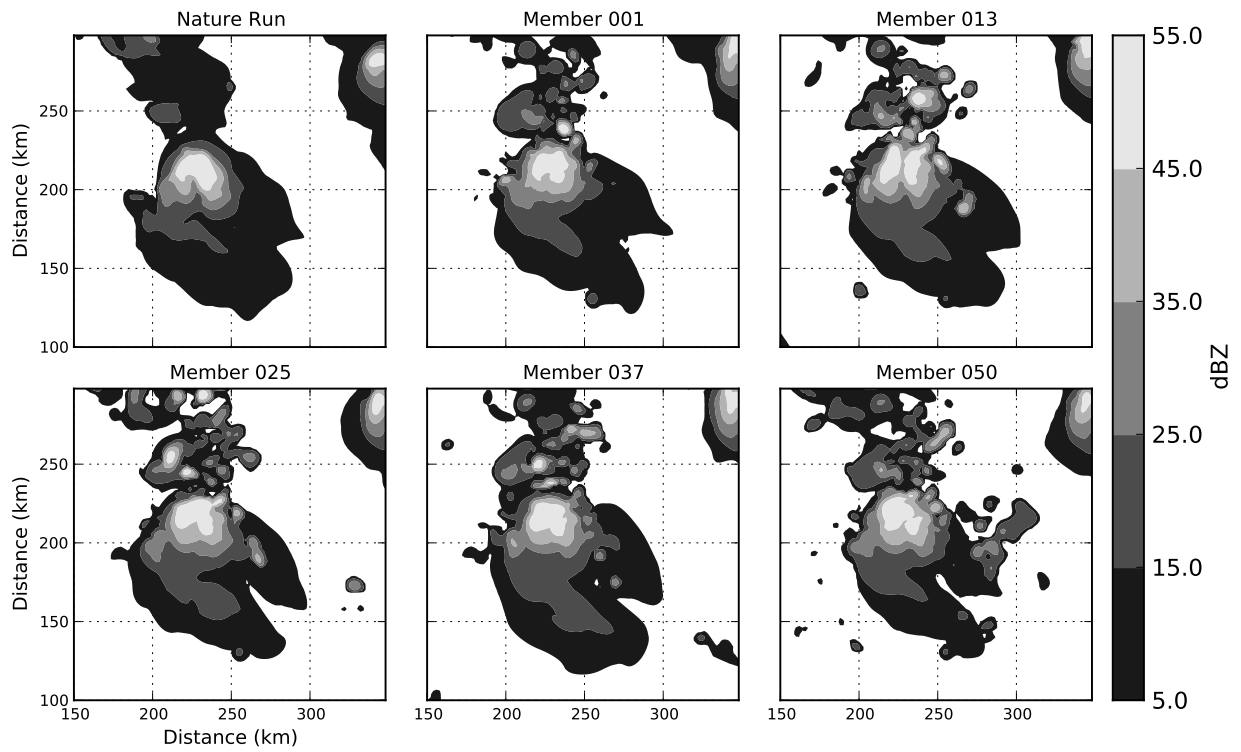


FIG. 8. Like Fig. 7, for R16.

Distribution of Analysis Ensembles (R4 and R16) around Nature Run (Realization 01) at 17 UTC

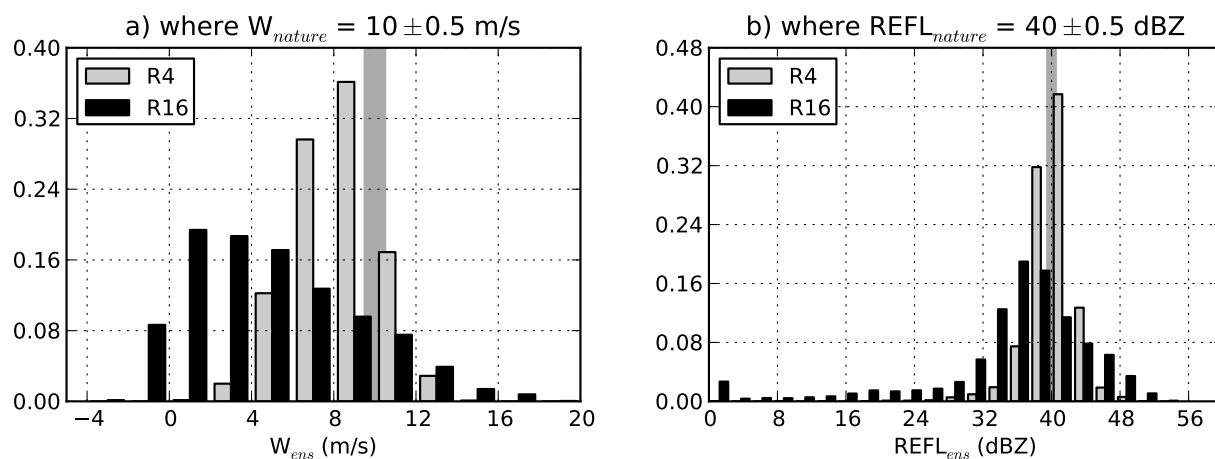


FIG. 9. Relative frequencies of model values within the ensemble-members of R4 (black bars) and R16 (gray bars). Regions inside storms are given by a) and b): a) at gridpoints where  $W$  of the nature run is  $10 \pm 0.5$  m s<sup>-1</sup>, b) where  $REFL_{nature} = 40 \pm 0.5$  dBZ. Regions closer than 64 km to the lateral boundaries are excluded so the histograms are computed only where the filter is fully active in R4 and R16. The bar width is 1/2 of the binwidth.

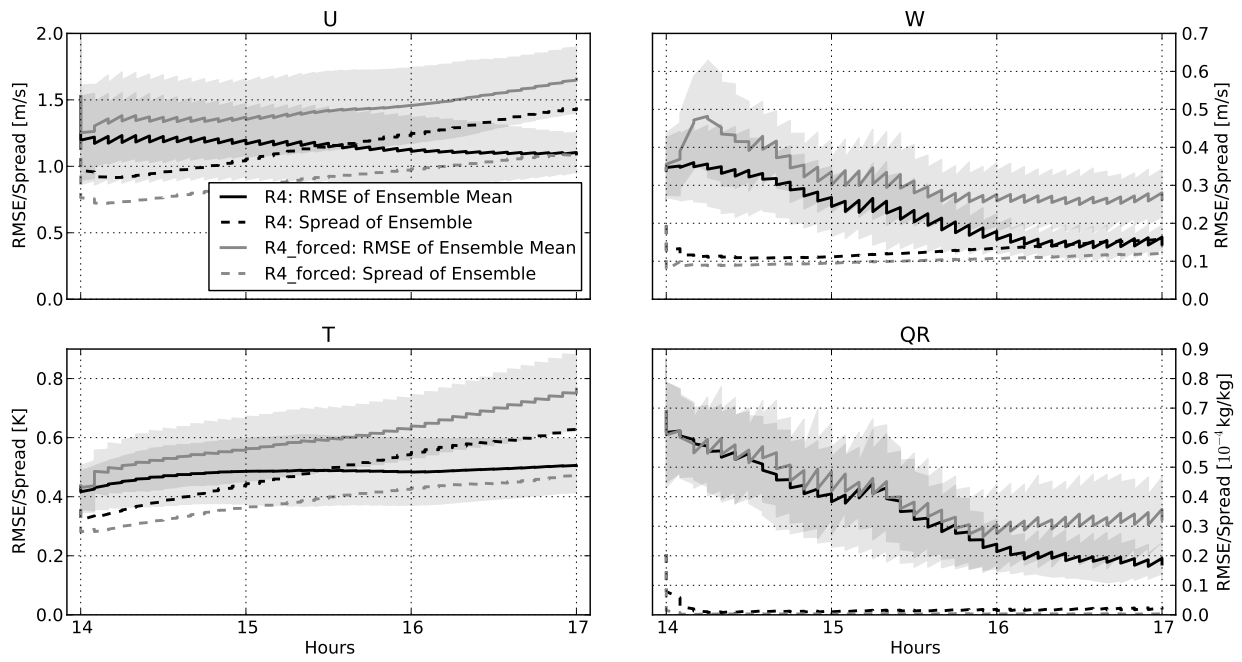


FIG. 10. Like Fig. 5, but comparing R4 (black) to R4\_forced (gray) for the primary variables  $U$ ,  $W$ ,  $T$  and  $QR$  in the first two hours of assimilation.

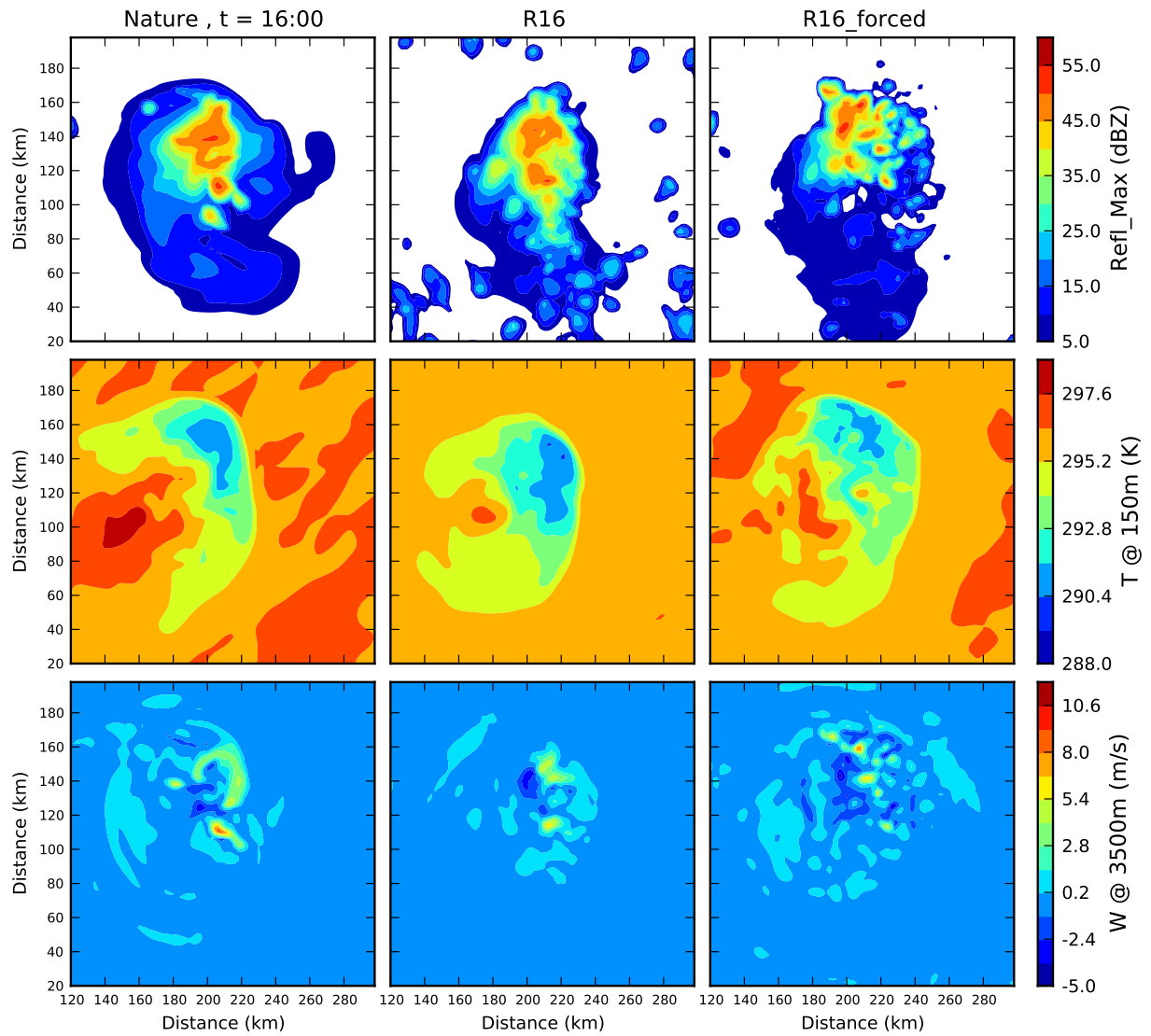


FIG. 11. Like Fig. 4, but comparing the analysis means of R16 and R16\_forced to the nature run at 16 UTC during the assimilation (Realization 03).

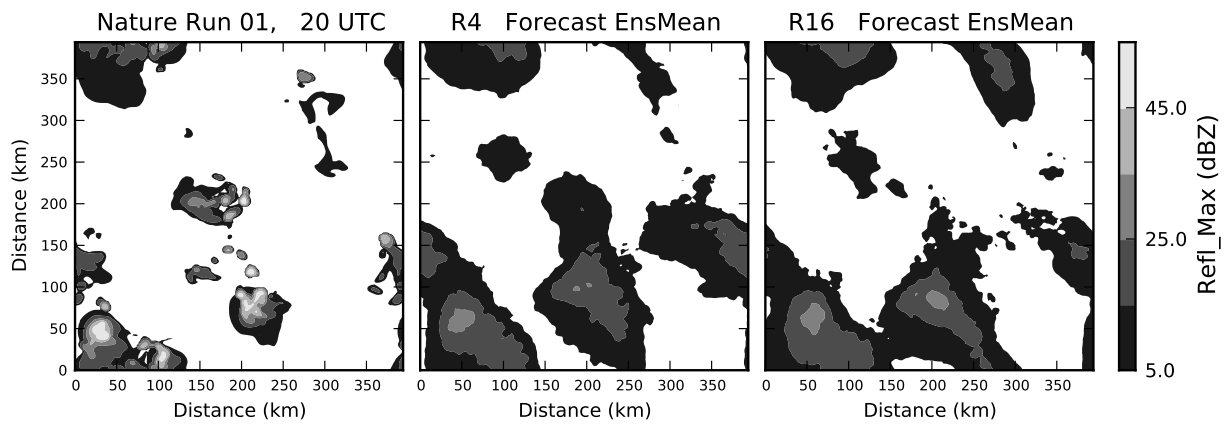


FIG. 12. Like Fig. 4, but forecast ensemble means (Realization 01) at 20 UTC.



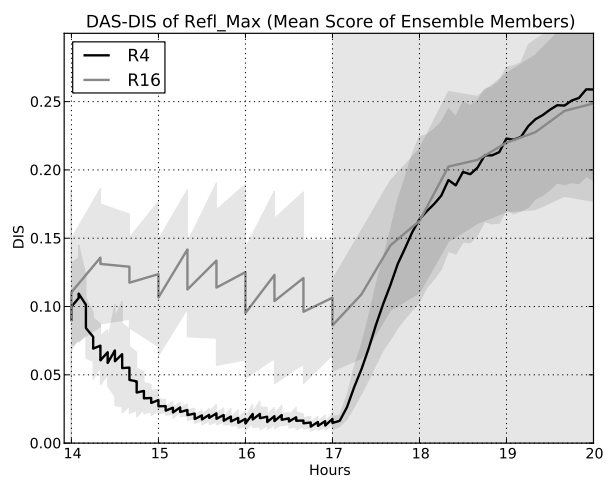


FIG. 13. DIS-component of the DAS-score of the maximum column reflectivity  $REFL\_MAX > 10$  dBZ. Displayed is the mean DIS-score of the ensemble members of R4 (black), R16 (gray). The error lines are the average over five experimental repetitions. An envelope of the empirical standard deviation is given by the shaded areas. A value of  $DIS = 0$  is a perfect match.

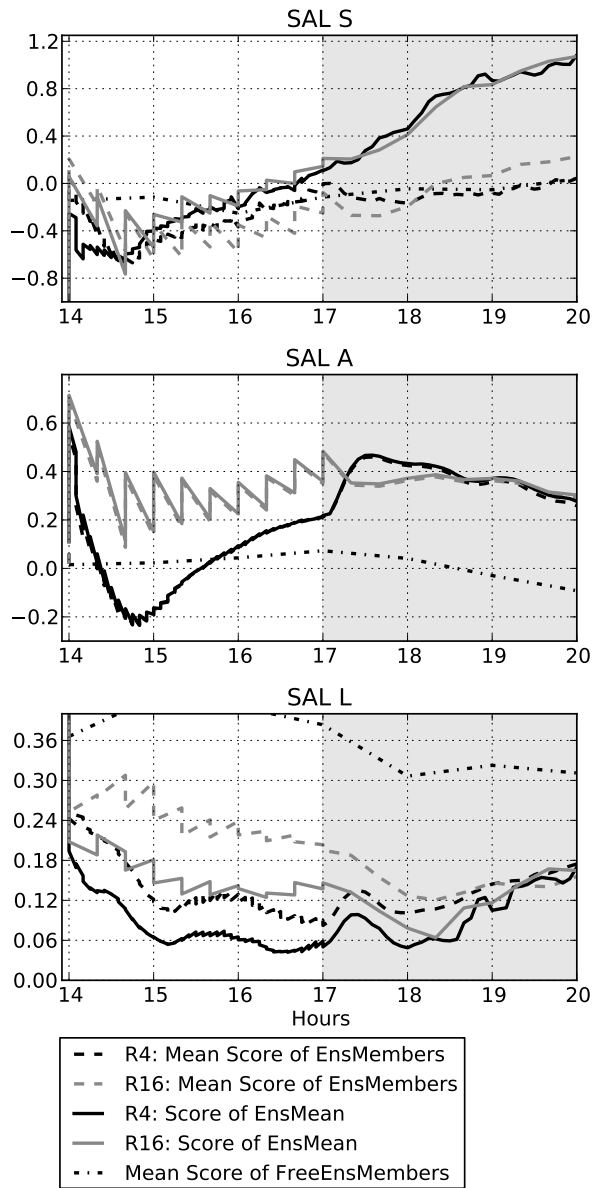


FIG. 14. SAL score components S, A and L (see text) of  $REFL\_MAX > 10$  dBZ, averaged over the five repetitions.

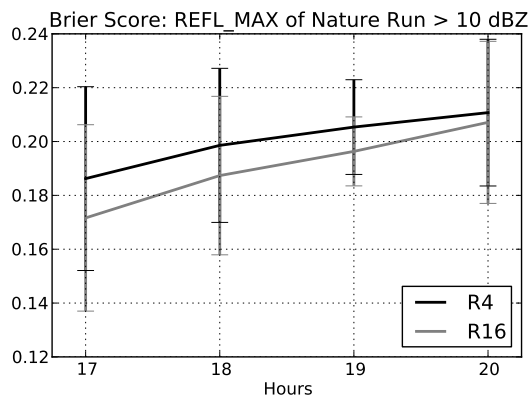


FIG. 15. Brier Score of forecast ensembles R4 and R16 at locations where  $REFL\_MAX > 10$  dBZ in the nature run. Errorbars show the empirical standard deviation of the five experimental repetitions.