

1 **Comments on: Thermodynamic characteristics of downdrafts in**  
2 **tropical cyclones as seen in idealized simulations of different**  
3 **intensities, by J. B. Wadler, D. S. Nolan, Jun A. Zhang and Lynn**

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7 In their interesting paper, [Wadler et al. \(2021a\)](#) examine the thermodynamic effect of  
8 downdrafts on the boundary layer and nearby updrafts in idealized simulations of category-3  
9 and category-5 tropical cyclones. These simulations are designated Ideal3 and Ideal5, re-  
10 spectively. The authors show that in the stronger storm, Ideal5, “downdrafts underneath  
11 the eyewall pose no negative thermodynamic influence because of eye-eyewall mixing below  
12 2-km altitude. Additionally, a layer of higher  $\theta_e$  between 1 and 2 km altitude associated  
13 with low-level outflow that extends 40 km outward from the eyewall region creates a “ther-  
14 modynamic shield” that prevents negative effects from downdrafts”. In the weaker storm,  
15 Ideal3, no such outflow occurs. The impact of downdrafts beyond the main eyewall in the  
16 two simulations are compared in a cartoon schematic, their Fig. 17.

17 In the case of Ideal5, a lower-tropospheric radial outflow jet is found to advect enhanced  $\theta_e$   
18 air outwards from the interior eyewall/eye region. Presumably,  $\theta_e$  is the pseudo-equivalent  
19 potential temperature. The radial outflow jet of enhanced  $\theta_e$  air is argued to act as a  
20 thermodynamic shield that modifies downdraft air descending into the boundary layer. As  
21 a downdraft attempts to pass through the shield, downdraft air would be warmed and  
22 moistened through mixing with the air in the shield. The net effect of the shield is to elevate  
23 the downdraft  $\theta_e$  before it is drawn inwards by the frictional inflow. In contrast, for the case  
24 of Ideal3, there is no appreciable persistent outflow jet to shield the vortex from downdraft  
25 influences.

26 In the conclusions they write: “The presence of a high- $\theta_e$  air above the inflow layer in  
27 Ideal5, which is also discussed in previous observational studies (e.g., [Barnes 2008](#); [Wadler  
28 et al. 2021b](#)), highlights the importance of storm structure in determining the thermodynamic  
29 effect of downdrafts. However, it remains unknown the exact mechanisms which lead to the  
30 formation of the high- $\theta_e$  above the boundary layer in the TC and why this layer formed in  
31 Ideal5, but not in Ideal3. This will be a topic of future work.”

32 In this comment, we offer an explanation for this feature, which we believe to be a common  
33 feature of vortex evolution as a storm matures and decays, especially in more intense storms

34 with a broad tangential wind circulation. The reasons are discussed in papers by [Kilroy](#)  
35 [et al. \(2016\)](#) and [Smith et al. \(2021\)](#). In essence, the evolution of a tropical cyclone at a  
36 particular stage of its life cycle depends broadly on the rate at which moist air is funnelled  
37 by the boundary-layer inflow towards the inner region, where deep convection prevails, and  
38 the rate at which this mass can be carried to the upper troposphere by the aggregate effects  
39 of this convection.

40 Typically, in the early stages of tropical cyclone formation and intensification, the bound-  
41 ary layer inflow is relatively weak and deep convection is more than capable of removing mass  
42 at the rate at which it is funnelled inwards. However, as the storm intensifies and the tan-  
43 gential wind field expands, the rate at which air is funnelled inwards increases, while the  
44 progressive warming of the upper troposphere tends to reduce the degree of convective insta-  
45 bility and thereby the ability of inner-core deep convection to ventilate the mass converging  
46 in the boundary layer. This reduced convective instability may be accompanied by the  
47 reduction of effective buoyancy in the eyewall as the eyewall broadens in size ([Smith and](#)  
48 [Montgomery 2022](#)). As soon as the boundary layer inflow begins to dominate, the residual  
49 mass that cannot be ventilated by deep convection flows outwards in a shallow layer just  
50 above the boundary layer. Typically, the tangential velocity component of air ascending out  
51 of the boundary layer in the inner-core region is supergradient ([Smith and Vogl 2008](#); [Smith](#)  
52 [et al. 2008, 2009](#) and has a natural tendency to flow outwards. Unless this air reaches a level  
53 of free convection, it remains stably stratified which accounts for the outflow occurring in a  
54 shallow layer.

55 [Smith et al. \(2021\)](#) showed that as a tropical cyclone matures, the low-level radial outflow  
56 becomes more and more prevalent and leads ultimately to the decay of the vortex, even in  
57 a quiescent environment. At radii where the radial outflow is sufficiently large so that the  
58 radial advection of absolute angular momentum,  $M$ , exceeds the vertical advection of  $M$  out  
59 of the boundary layer, the tangential flow at the top of the boundary layer will spin down.  
60 The foregoing evolution is illustrated, for example, in Figs. 4 and 5 of [Smith et al. \(2021\)](#).

61 The difference between the two simulations Ideal3 and Ideal5 is presumably because, at  
62 the time of analysis, Ideal3 is not yet in a state where the mass influx in the boundary layer  
63 exceeds the rate at which this flux can be ventilated by deep convection. It might be possible  
64 to validate this conjecture using the ventilation diagnostic introduced by [Smith et al. \(2021\)](#):  
65 see their Eq. (10) and Fig. 8.

66 Axisymmetric balance models of tropical cyclone evolution can mimic the same effect,  
67 even though the winds in the boundary layer do not become supergradient in such a model  
68 by definition. This was shown in a series of simulations with a prognostic balance model  
69 in which the diabatic heating rate is varied in strength in relation to the frictional forcing  
70 ([Smith and Wang \(2018\)](#)). In this model, the unbalanced forces do not exist.

71 On a separate topic, it is worth noting that the [Wadler et al.](#) study (p3517) demonstrates  
72 the quantitative importance of “low-level outflow from inside the eye and eye-eyewall mixing”  
73 in supporting relatively high  $\theta_e$  air ascending the eyewall in their Ideal5 vortex. The “eye-  
74 eyewall mixing” of near surface boundary layer  $\theta_e$  was invoked by [Persing and Montgomery](#)  
75 ([2003](#)) to explain how hurricanes could significantly exceed Emanuel’s Potential Intensity  
76 (PI) theory for the maximum gradient wind. However, [Bryan and Rotunno \(2009\)](#) showed  
77 this process was too weak to explain the discrepancy between PI theory and numerical  
78 experiments in strictly axisymmetric model configurations. As discussed in [Montgomery](#)  
79 [and Smith \(2017\)](#), the dynamical origin of “super-intense storms” is now better understood.  
80 In particular, there is an important distinction between the radius of maximum gradient  
81 wind  $r_{gm}$  and the radius of maximum tangential wind  $r_m$ , the former typically lying some 10  
82 - 20 km (or more) outside the latter. Scientific questions remain concerning the impact of  
83 “high-octane”  $\theta_e$  air generated at radii inside  $r_{gm}$  (including the low-level eye of the storm)  
84 in supporting locally buoyant updrafts comprising a realistic, three-dimensional eyewall.  
85 The [Wadler et al.](#) results re-affirm the non-negligible influence of eye-eyewall mixing in a  
86 three-dimensional, category-5, hurricane.

87     *Acknowledgments.*

88     MTM acknowledges the support of NSF AGS-1313948, NOAA HFIP grant N0017315WR00048,  
89    NASA grant NNG11PK021 and the U.S. Naval Postgraduate School.

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